Quasi-Linear Differential Equations Two Second Order Equations

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Symbiont Workshop, March 15th 2022, Bonn

Objectives

Existence, (Non)Uniqueness and Regularity of Solutions of an Initial Value Problem

 $u(\overline{x}) = \overline{u}, \quad u'(\overline{x}) = \overline{u}'$

for a Second Order ODE of Type

g(x)u'' = f(x, u, u')

with \overline{x} a zero of g(x).

Two Cases: Initial condition \overline{x} is a simple or a double zero of g(x)

Why relevant for Symbiont?

Connection between quasi-linear systems $A(\mathbf{u})\mathbf{u}' = F(\mathbf{u})$ and singular pertubation problems $\mathbf{u}' = f(\mathbf{u}, \mathbf{v}, \epsilon), \ \epsilon \mathbf{v}' = g(\mathbf{u}, \mathbf{v}, \epsilon)$

An Analytic View

Theorem (Liang 2009) Let g(x) = x and set $\gamma = f_{u'}(0, \overline{u}, \overline{u'}), k = \lceil \gamma \rceil$ $\gamma < 0 \Rightarrow$ unique solution $u \in C^{\infty}([0, \delta])$ for some $\delta > 0$. $\gamma > 0 \Rightarrow \exists \overline{u}^{(2)}, \dots, \overline{u}^{(k)} \in \mathbb{R}$ (uniquely determined by $\overline{u}, \overline{u'}, f$): for every solution

$$\lim_{x \to 0+} \frac{1}{x^{\gamma}} \left(u'(x) - \sum_{j=1}^{k} \frac{\overline{u}^{(k)} x^{j-1}}{(j-1)!} \right)$$

exists, each solution uniquely determined by limit; one solution smooth, all others in $C^k \setminus C^{k+1}$.

• $\gamma \in \mathbb{N} \Rightarrow$ additional logarithmic term and either all solutions are smooth or all in $C^k \setminus C^{k+1}$ ("resonance").

Theorem (Brunovsky, Cerny, Winkler 2013)

$$x^2u'' = axu' + bu - c(u'-1)^2, \quad x \in (0, x_0)$$

with $a, b \in \mathbb{R}$, c > 0 and initial condition u(0) = 0.

- a + b < 0 ⇒ no continuous solutions exist</p>
- a + b > 0 ⇒ infinitely many continuous solutions exist; unique solution u corresp. to the choice x₀ = ∞ with 0 ≤ u(x) ≤ x for all x > 0 and u is strictly increasing and concave.

The Geometry of ODEs

- 1. Jet bundle \mathcal{J}_q the set of all equivalence classes $[\phi]_x^{(q)}$ of smooth functions $\phi: \mathbb{R} \to \mathbb{R}$ with the same Taylor expansion at x up to order q
- 2. Coordinates $(x, u, u', \ldots, u^{(q)})$
- 3. Natural projection $\pi_r^q : \mathcal{J}_q \to \mathcal{J}_r$ for q > r and $\pi^q : \mathcal{J} \to \mathbb{R}$.
- 4. The contact distribution $C^{(q)}$: It is spanned by the two vector fields

$$C^{(q)} = \partial_x + u^{(1)}\partial_u + \dots + u^{(q)}\partial_{u^{(q-1)}}$$
 (transversal)
and $C_q = \partial_{u^{(q)}}$ (vertical)

- 5. Identify a scalar *q*th-order differential equation F = 0 as a submanifold $\mathcal{R}_q \subseteq \mathcal{J}_q$
- Vessiot space V_ρ[R_q] = T_ρR_q ∩ C^(q)|_ρ at ρ ∈ R_q. Computation: Every X ∈ V_ρ[R_q] has a representation as

$$X = aC^{(q)}|_{\rho} + bC_q|_{\rho}$$

where a, b are the solutions of

$$C^{(q)}(F)(\rho)a + C_q(F)(\rho)b = 0$$

Example Sphere

Consider the equation

$$F = u^{\prime 2} + u^2 + x^2 - 1 = 0$$

in \mathcal{J}_1 with contact disrtibution $C^{(1)} = \partial_x + u' \partial_u$ and $C_1 = \partial_{u'}$.



Example Sphere

Vessiot distribution:

$$aC^{(2)}(F) + bC_2(F) = a(2x + 2uu') + b2u' = 0$$
$$\rightarrow X = u'\partial_x + u'^2\partial_u - (x + uu')\partial_{u'}$$



Solutions

- 1. Identify a function $\phi : \mathbb{R} \to \mathbb{R}$ with its section $\sigma : \mathbb{R} \to \mathbb{R}^2$ where $\sigma(x) = (x, \phi(x))$.
- 2. Prolonged section

$$j_q \sigma_\phi : \mathbb{R} \to \mathcal{J}_q, \quad x \mapsto (x, \phi(x), \phi'(x), \dots, \phi^{(q)}(x)).$$

- 3. A function ϕ is a strong solution of \mathcal{R}_q if $j_q \sigma_{\phi} \subseteq \mathcal{R}_q$
- 4. A generalised solution is a one dimensional integral manifold $\mathcal{N} \subseteq \mathcal{R}_q$ of $\mathcal{V}[\mathcal{R}_q]$.
- 5. A generalised solution is **proper** if there does not exist x such that $\mathcal{N} \subseteq (\pi^q)^{-1}(x)$.
- 6. The projection $\pi_0^q(\mathcal{N}) \subset \mathcal{J}_0$ of a proper generalised solution is called a **geometric** solution.



Singularities

Two Types of Singularities:



Definition

Let ρ be a smooth point of \mathcal{R}_q .

- 1. If $\dim(\mathcal{V}_{\rho}[\mathcal{R}_q]) > 1$, then ρ is called an irregular singularity.
- 2. If $\dim(\mathcal{V}_{\rho}[\mathcal{R}_q]) = 1$ and $\mathcal{V}_{\rho}[\mathcal{R}_q]$ is vertical, then ρ is called a regular singularity.
- 3. Otherwise ρ is called regular.

Examples Singularities and Solutions





No Irregular Singularities

Theorem

Let $\mathcal{R}_q \subset \mathcal{J}_q$ be a scalar ordinary differential equation such that at every $\rho \in \mathcal{R}_q$ the Vessiot space $V_\rho[\mathcal{R}_q]$ is one-dimensional.

- 1. If ρ is regular, then \exists_1 strong (two sided) solution σ with $\rho \in im(j_q \sigma)$.
- If ρ is regular singular, then either two strong one sided solutions σ₁, σ₂ with σ_i ∈ imj_qσ_i exist or only one strong two-sided solution exists whose (q + 1)th derivative blows up at x = π^q(ρ).



Prolongation

Recall: \mathcal{R}_q is defined by the zero set of the function $F : \mathcal{J}_q \to \mathbb{R}$. The first prolongation $\mathcal{R}_{q+1} \subseteq \mathcal{J}_{q+1}$ is the zero set defined by F and the function

$$D_{\mathsf{x}}\mathsf{F} = C^{(q)}(\mathsf{F}) + C_q(\mathsf{F})u^{(q+1)} : \mathcal{J}_{q+1} \to \mathbb{R}$$

Proposition

Let $\mathcal{F}_{q+1} = (\pi_q^{q+1})(\rho_q) \cap \mathcal{R}_{q+1}$ be the fibre of a point $\rho_q \in \mathcal{R}_q$ in the first prolongation \mathcal{R}_{q+1} .

- 1. If ρ_q is regular, then $\mathcal{F}_{q+1} \neq \emptyset$ and consists of regular points of \mathcal{R}_{q+1} .
- 2. If ρ_q is regular singular, then $\mathcal{F}_{q+1} = \emptyset$.
- 3. If ρ_q is irregular singular, then $\mathcal{F}_{q+1} \neq \emptyset$ and consists of singular points of \mathcal{R}_{q+1} .

Quasi-Linear Equations

We consider now quasi-linear equations type

$$g(x, u_{q-1})u^{(q)} = f(x, u_{q-1}).$$

Equation determining the Vessiot space at a point $\rho = (x, \overline{u}_{q-1}) \in \mathcal{R}_q$ becomes

$$[C^{(q)}(g)(\rho)\overline{u}^{(q)} - C^{(q)}(f)(\rho)]a + g(\rho)b = 0$$

 \Rightarrow a point ρ is singular if $g(\rho) = 0$ (independent of $\overline{u}^{(q)}$) and it is irregular singular if and only if in addition

$$C^{(q)}(g)(\rho)\overline{u}^{(q)} - C^{(q)}(f)(\rho) = 0.$$

The Vessiot space is

$$X = g(x, \boldsymbol{u}_{q-1})C^{(q)} - [C^{(q)}(g(x, \boldsymbol{u}_{q-1}))u^{(q)} - C^{(q)}(f(x, \boldsymbol{u}_{q-1}))]C_q$$

Key Property

Key Property of Quasi-linear Differential Equations: X is π_{q-1}^q -projectable to

$$Y = g(x, u_{q-1})C^{(q-1)} + f(x, u_{q-1})C_{q-1}.$$

Assumption: Y is continueable to all of \mathcal{J}_{q-1} .

Definition

- 1. A point $\tilde{\rho}_{q-1} \in \mathcal{J}_{q-1}$ is called impasse point for $\mathcal{R}_q \subseteq \mathcal{J}_q \iff Y$ is not transversal to $\pi^{(q-1)}$ at $\tilde{\rho}_{q-1}$ (∂_x -component vanishes)
- 2. An impasse point $\tilde{\rho}_{q-1} \in \mathcal{J}_{q-1}$ is called **proper** \iff Y vanishes at $\tilde{\rho}_{q-1}$. Otherwise it is called improper.

The local analysis of proper impasse points reduces to the study of the stationary points of the vector field Y.

Impasse Points of Quasi-Linear Equations

Definition

- 1. A one-dimensional invariant manifold of Y is called a weak generalised solution.
- The π^{q-1}₀-projection of a weak generalised solution is called a weak geometric solution.

No guarantee that qth derivative exists \rightsquigarrow generally not related to a classical solution

First Type of Quasi-Linear Equations

Problem 1:

For a quasi-linear differential equation of type

$$g(x)u^{\prime\prime}=f(x,u,u^{\prime})$$

we consider the singular initial value problem $\rho_1 = (\overline{x}, \overline{u}_0, \overline{u}_1)$ where $f(\rho_1) = 0$ and \overline{x} is a simple zero of g(x).

Prolonged Equations and Vessiot Distribution

For q>2 compute recursively the prolonged equations \mathcal{R}_q defined by F_2,\ldots,F_{q-1} and

$$F_q(x, \boldsymbol{u}_{(q)}) = g(x)u^{(q)} + [(q-2)g'(x) - f_{u'}(x, u, u')]u^{(q-1)} - h_q(x, \boldsymbol{u}_{(q-2)})$$

where

$$h_3(x, \boldsymbol{u}_{(1)}) = C^{(1)}(f(x, u, u'))$$

$$h_q(x, \boldsymbol{u}_{(q-2)}) = C^{(q-2)}(h_{q-1}(x, \boldsymbol{u}_{(q-3)}) - [(q-3)g'(x) - f_{u'}(x, u, u')]u^{(q-2)})$$

On the prolonged equations \mathcal{R}_q the Vessiot distribution is given by

$$X^{(q)} = g(x)C^{(q)} + \left([(q-1)g'(x) - f_{u'}(x, u, u')]u^{(q)} - h_{q+1}(x, u_{q-1}) \right) C_q$$

Singularities: $\rho_q = (\overline{x}, \overline{u}, \overline{u}', ..., \overline{u}^{(q)}) \in \mathcal{R}_q$

1. regular singularity $\iff g(\overline{x}) = 0$

2. irregular singularity $\iff g(\overline{x}) = 0$ and $[g'(\overline{x}) - f_{u'}(\rho_1)]\overline{u}^{(q)} = h_{q+1}(\rho_{q-1})$ where $\pi_i^q(\rho_q) = \rho_i$ for i = q - 1 and i = 1.

The Dynamical System

The Projected Vessiot distribution is

$$Y^{(q-1)} = g(x)C^{(q-1)} + (h_q(x, \boldsymbol{u}_{(q-2)}) - [(q-2)g'(x) - f_{u'}(x, \boldsymbol{u}_{(1)})]u^{(q-1)})C_{q-1}$$

For our singular initial values $\rho_1 = (\overline{x}, \overline{u}, \overline{u}')$ all points

$$\rho_q = (\overline{x}, \overline{u}_q) \in \mathcal{R}_q \cap (\pi_1^q)^{-1}(\rho_1)$$

are singular too and the projection $\rho_{q-1} = \pi_{q-1}^q(\rho_q) \in \mathcal{R}_{q-1}$ is a proper impasse point of \mathcal{R}_q

 \rightsquigarrow analyse the local phase portrait of $Y^{(q-1)}$ at ho_{q-1}

The Jacobian at the Singularity

We set $\delta = g'(\overline{x})$ and $\gamma = f_{u'}(\rho_1)$ and assume wlog that $\delta > 0$. At $\rho_{q-1} = (\overline{x}, \overline{u}_0, \dots, \overline{u}_{q-1})$ the Jacobian of $Y^{(q-1)}$ is

$$J^{(q-1)} = \begin{pmatrix} \delta & 0 & \cdots & 0 \\ \delta \overline{u}_1 & 0 & \cdots & 0 \\ \vdots & \vdots & & \vdots \\ \delta \overline{u}_{q-1} & 0 & \cdots & 0 \\ a_0 & \cdots & a_{q-1} & \gamma - (q-2)\delta \end{pmatrix}$$

with eigenvalues δ , 0 ((q-1) times) and $\gamma - (q-2)\delta$.

 \mathcal{R}_{q-1} is a 3-dimensional manifold \rightsquigarrow

 δ , 0 and $\gamma - (q-2)\delta$ are the relevant eigenvalues

Resonance

Definition

We say that the initial value problem at ρ_1 has a **resonance** at order $k \in \mathbb{N}$, if

 $k\delta = \gamma.$

In this case at any point $\rho_k \in (\pi_1^k)^{-1}(\rho_1)$ above ρ_1 we call $A_k = h_{k+2}(\rho_k)$ the resonance parameter. The resonance is critical at ρ_k if $A_k \neq 0$ and smooth if $A_k = 0$.

Proposition

Let $\rho_q \in \mathcal{R}_q$ be an irregular singularity and let $\mathcal{F}_{q+1} = (\pi_q^{(q+1)})^{-1}(\rho_q)$ be the fibre.

- 1. Then $\mathcal{F}_{q+1} \subset \mathcal{R}_{q+1}$
- 2. If ρ_1 is not in resonance at order q, then \mathcal{F}_{q+1} contains exactly one irregular singularity.

3. If ρ_1 is in $\left\{\begin{array}{c} critical\\ smooth \end{array}\right\}$ resonance at order q, then \mathcal{F}_{q+1} consists entirely of $\left\{\begin{array}{c} regular\\ irregular \end{array}\right\}$ singularity.

Eigenvectors

Eigenvectors of $J^{(q-1)}$:

Without resonance:

Resonance at order q - 1: (Generalised) Eigenspace generated by

$$e_1 = (0, \dots, 0, 1)^T$$
 and $e_2 = (1, \overline{u}', \dots, \overline{u}^{(q-1)}, 0)^T$

Resonance smooth: e₁ and e₂ are proper eigenvectors

Resonance critical: e₁ proper and e₂ generalised eigenvector

The Case Without Resonance

Theorem

Assume that at no order a resonance appears.

- 1. If $\delta\gamma < 0$, then the corresponding initial value problem possesses a unique two-sided smooth solution and no additional one-sided solutions.
- If δγ > 0, then there exists a one-parameter family of two-sided solutions. One of these solutions is smooth; the other ones are in C^k \ C^{k+1} with k = [γ/δ]. All of these solutions possess the same Taylor polynomial ∑^k_{i=0} ^{ū_i}/_{i!}(x − x̄)ⁱ of degree k around x̄.

The Case of $\gamma = 0$

Theorem

Assume that $\gamma = 0$. Then there exists a unique smooth two-sided solution (and possibly further one-sided solutions).

Theorem

Assume that a resonance occurs at the order k > 0. There exists a one-parameter family of two-sided solutions all possessing the same Taylor polynomial $\sum_{i=0}^{k} \frac{\overline{u}_i}{i!} (x - \overline{x})^i$ of degree k around \overline{x} . In the case of a smooth resonance, all of these solutions are smooth and each is uniquely determined by the value of its (k + 1)st derivative in \overline{x} . In the case of a critical resonance, all solutions live in $C^k \setminus C^{k+1}$.

Brunovsky

Problem 2: For the quasi-linear differential equation

$$x^2u'' = axu' + bu - c(u' - 1)^2$$

we consider the singular initial value problem $\rho_1 = (\overline{x}, \overline{u}_0, \overline{u}_1) = (0, 0, 1)$ with are real parameters a, b, $c \in \mathbb{R}$.

Joint work with Peter Szmolyan

The projected Vessiot distribution is

$$Y = x^2 \partial_x + x^2 u' \partial_u + (axu' + bu - c(v-1)^2) \partial_{u'}.$$

The set of stationary points is given by the parabola

$$u=\frac{c}{b}(u'-1)^2$$

and the Jacobian has at ρ_1 triple eigenvalue zero.

 \Rightarrow desingularise the stationary point ρ_1 with a blow-up

Consider the transformed system

$$(b(ay_1y_3 - cy_1^2) + y_2)\partial_{y_1} + (by_3^2(b^2y_1 + a + b))\partial_{y_2} + by_3^2\partial_{y_3}$$

with Jacobian

$$\begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

at the origin.

Blow-up in positive and negative y_1 direction:

Stationary points:

$$s^{1} = (0, 0, 0), s^{2} = (0, bc, 0) \text{ (later)}$$

 $J(s^{1}) = \begin{pmatrix} -bc/2 & 0 & 0 \\ 0 & 2bc & 0 \\ 0 & 0 & 3bc/2 \end{pmatrix}$

 s^1 repelling/attracting node y_1 direction attracting/repelling



Stationary points:

$$s^{3} = (0, bc, 0), s^{4} = (0, 0, 0) \text{ (later)}$$

 $J(s^{4}) = \begin{pmatrix} bc/2 & 0 & 0 \\ 0 & -2bc & 0 \\ 0 & 0 & -3bc/2 \end{pmatrix}$

 s^4 attracting/repelling node y_1 direction repelling/attracting





Blow up in positive y_2 direction (negative analogously):

Stationary points

$$s^2=(lpha,0,0)$$
 and $s^3=(-lpha,0,0)$ with $lpha=\sqrt{rac{1}{bc}}$

with Jacobians

$$J(s^2) = \begin{pmatrix} -2bc\alpha & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \text{and} \quad J(s^3) = \begin{pmatrix} 2bc\alpha & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

Lemma

- 1. Let $a + b \neq 0$ and bc < 0. Then in the chart for the positive y_2 -direction there are no real stationary points on the plane $\{y_2 = 0\}$.
- Let a + b ≠ 0 and bc > 0. Then in the chart for the positive y₂-direction we find on the plane {y₂ = 0} the two stationary points

$$s^2=(lpha,0,0)$$
 and $s^3=(-lpha,0,0).$

On the plane { $y_2 = 0$ } they show the following dynamical behaviour: 2.1 For b(a + b) > 0 the point s^2 is an attracting node and s^3 is a saddle. 2.2 For b(a + b) < 0 the point s^2 is a saddle and s^3 is a repelling node.



At s^2 and s^3 the Jacobians have double eigenvalue zero

 \Rightarrow Compute center manifolds and their dynamics

Lemma

On the center manifold $C(s^2)/C(s^3)$ the point s^2/s^3 is a saddle point. On the halfspace $\{y_2 \ge 0\}$ we have the following dynamics on $C(s^2)/C(s^3)$:

- 1. a + b > 0, b > 0: attracting directions $\pm (0, 0, 1)^{tr}$ and repelling direction $(0, a + b, 1)^{tr}$.
- 2. a + b > 0, b < 0: repelling directions $\pm (0, 0, 1)^{tr}$ and attracting direction $(0, a + b, 1)^{tr}$.
- 3. a + b < 0, b > 0: repelling directions $\pm (0, 0, 1)^{tr}$ and attracting direction $-(0, a + b, 1)^{tr}$.
- 4. a + b < 0, b < 0: attracting directions $\pm (0, 0, 1)^{tr}$ and repelling direction $-(0, a + b, 1)^{tr}$.

- 1. a + b > 0, b > 0: attracting directions $\pm (0, 0, 1)^{tr}$ and repelling direction $(0, a + b, 1)^{tr}$.
- 2. a + b > 0, b < 0: repelling directions $\pm (0, 0, 1)^{tr}$ and attracting direction $(0, a + b, 1)^{tr}$.
- 3. a + b < 0, b > 0: repelling directions $\pm (0, 0, 1)^{tr}$ and attracting direction $-(0, a + b, 1)^{tr}$.
- 4. a + b < 0, b < 0: attracting directions $\pm (0, 0, 1)^{tr}$ and repelling direction $-(0, a + b, 1)^{tr}$.





Putting it together



 \Rightarrow the stationary point is a saddle node